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Brajesh Tiwari, Rajveer Jha, and V. P. S. Awana

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Magnetotransport studies of FeSe under hydrostatic pressure

Brajesh Tiwari, Rajveer Jha, and V. P. S. Awana^a

CSIR-National Physical Laboratory, Dr. K.S. Krishnan Marg, New Delhi-110012, India

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The discoveries of iron-based superconductors with relatively high transition temperature are under intense experimental and theoretical investigation. Here we present magnetotransport measurements on FeSe superconductor under hydrostatic pressure. We show that in Fe-deficient tetragonal FeSe binary compound, the onset of superconducting transition is almost doubled under 1.98 GPa pressure and the estimated upper critical field of 26.7 Tesla is increased to 47.5 Tesla. © 2014 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution 3.0 Unported License. [<http://dx.doi.org/10.1063/1.4885473>]

I. INTRODUCTION

Superconductivity in iron-based compounds is fast developing research subject in condensed matter physics and materials science community.¹⁻⁴ Since most of the Fe-based compounds show magnetic ordering at low temperatures, the unexpected discovery of 26 K superconducting transition in $\text{LaFeAsO}_{1-x}\text{F}_x$ by Kamihara *et al.*¹ in 2008 rejuvenate the superconducting research community. Fe-based superconductors share common features with their relatives' high-T_c Cu-based (curates) superconductors for example, both are d-electron materials having layered structure, their parent compounds are antiferromagnetic and superconductivity arises from electron or/and hole doping or by application of pressure. The detailed review on this topic can be seen in Ref 2 and Ref 3.^{2,3} A simple binary compound FeSe in tetragonal phase has emerged as a superconductor with T_c of ~ 8-13 K.⁴⁻⁷ Essentially FeSe has identical packing to that of family of layered FeAs-based high-T_c superconductors however lacking intercalated charge reservoir layers.⁸ Iron arsenide and iron selenide share common structure which contains Fe_2X_2 (X = As and Se) layers of edge-shared FeX_4 tetrahedra suggesting same mechanism responsible for superconductivity in FeSe. In contrast to FeSe, chemical substitutions require for LiFeAs ,⁹ BaFe_2As_2 ,¹⁰ LnFeAsO ^{1,2} to drive them to superconductivity regime from itinerant antiferromagnetic state. Superconductivity in FeSe is very sensitive to stoichiometry, defects and disorder in the system.¹¹ In iron selenide, using various cationic spacer layers consisting of metal ions (for example Li^+ , Na^+ , K^+ , Ba^{2+} , Ca^{2+} , Eu^{2+} etc), the superconductivity could be increased to T_c of upto 45 K.^{12,13} Also using different molecular spacer layers, the superconducting transition is achieved as high as 56 K in FeSe based superconductors.^{14,15} The very first report on superconductivity of FeSe under pressure of 1.48 GPa was reported by Mizuguchi *et al.*, with increase of T_c from 13.5 K to 27 K.¹⁶ The detailed pressure evolution of low temperature structural studies revealed an intimate link between superconducting properties and the Se-FeSe interlayer separation.¹⁷ The nuclear magnetic resonance (NMR) measurements by Imai *et al.* suggested link between spin fluctuations and superconductivity in FeSe, however the Mossbauer studies confirmed that the same cannot be attributed to suppression of spin-density wave, unlike arsenide superconductors.^{7,18} In the present work, we revisited the superconductivity in FeSe under hydrostatic pressure via magnetotransport measurements. These measurements were performed on Fe-deficient tetragonal FeSe sample with minor hexagonal (NiAs-type) impurity.

^aCorresponding Author: Dr. V. P. S. Awana, Principal Scientist, E-mail: awana@mail.nplindia.org, Ph. +91-11-45609357, Fax:+91-11-45609310, Homepage www.ftewebs.com/vpsawana/



II. EXPERIMENTAL DETAILS

Polycrystalline bulk samples were prepared from high purity (>99%) powder of iron and selenium mixed and ground in an agate mortar and pestle with nominal stoichiometry of FeSe. The grinding was carried out in argon filed glove box. The grinded powder was cold pressed into rectangular pellet with 100kg/cm² uniaxial stress. The pellets were sealed in an evacuated quartz tube (in 10⁻⁴ bar vacuum) and temperature was ramped in a box Furness to 750°C i.e., above boiling point of Se (685°C) at the rate of 2°C/min. The sample was kept at this temperature for 12 h before cooling to room temperature by rate of 2°C/min. The sample was re-grinded and pelletized before final sintering at 750°C for 12 h. Powder x-ray diffraction (XRD) data was recorded at room temperature in 2θ range of 5° to 60° in step of 0.02 using Rigaku made diffractometer with Cu K α ($\lambda = 1.54 \text{ \AA}$). Multiphase Rietveld refinement was carried out using Full-Prof software. Physical Property Measurements System (PPMS-14T, *Quantum Design*) is used to perform magnetotransport measurements under hydrostatic pressure, which is applied using HPC-33 Piston cell in addition to DC resistivity Option. Hydrostatic pressures were generated by a BeCu/NiCrAl clamped piston-cylinder cell under a fluid medium of Fluorinert in a 3 Teflon cell. Annealed Pt wires attached to gold-sputtered contact surfaces on the sample by silver epoxy in a standard four-wire configuration. The applied pressure was calibrated from the superconducting transition temperature of Pb.

III. RESULTS AND DISCUSSION

X-ray powder diffraction pattern is Rietveld refined together for tetragonal P4/nmm and hexagonal P6₃/mmc structures of FeSe using Full-Prof and is shown in Figure 1(a). The global fitness of XRD pattern is $\chi^2 = 3.21$. Nonsymmorphic tetragonal structure of FeSe with space group P4/nmm (#129) is best fit ted to the lattice parameters $a = b = 3.774(2) \text{ \AA}$ and $c = 5.520(6) \text{ \AA}$ along with hexagonal (NiAs type) P6₃/mmc (#196) as a minority phase (with lattice parameters $a = b = 3.626(5) \text{ \AA}$ and $c = 5.890(1) \text{ \AA}$), which is close to the previous reported values.^{6,8,11,18} Both tetragonal and hexagonal crystal structures are represented in Figure 1(b). In tetragonal phase, Fe and Se occupy the Wyckoff positions; 2a (3/4, 1/4, 0) with site symmetry -4m2 and 2c (1/4, 1/4, 0.275) with site symmetry 4mm, respectively. However in hexagonal phase Fe occupies Wyckoff position; 2a (0,0,0) with site symmetry -3m and Se; 2c(1/3, 2/3, 1/4) with site symmetry -6m2. Tetragonal phase of FeSe is formed by layers of FeSe whereas hexagonal phase does not have this type of structure. The best fit of XRD pattern of FeSe gives Bragg R-factor of 3.45 and 3.28 for tetragonal and 20% of secondary hexagonal phases, respectively and deduced stoichiometry Fe_{0.954±0.005}Se in its tetragonal phase. When we forced the software to strict stoichiometry of FeSe, the obtained Bragg-R factor and reduced χ^2 are higher compared to the composition Fe_{0.954±0.005}Se and hence confirms the Fe-deficient FeSe in its tetragonal phase. This Fe-deficient tetragonal phase of FeSe shows superconducting transition which is consistent to the recently reported compositions by Chen *et al.*¹⁹ but in contrary to some earlier reports.⁵⁻⁸

Resistivity (ρ) measured under different external hydrostatic pressure upto 1.98 GPa is displayed in Figure 2 for the studied FeSe. Normal state resistivity follows the metallic behavior for FeSe, as can be seen in Figure 2(a). Interestingly, an anomalous deviation under increasing pressure can be observed in resistive ty at above ~ 100 K. Based on NMR investigations Imai *et al.* suggested that at below 100 K temperature, antiferromagnetic spin fluctuations are strongly enhanced towards superconducting transition.⁷ Magnified part of ρ -T curve can be viewed as inset of Figure 2(a) from which onset of T_c is determined by extrapolating the curve to straight lines (solid lines). The obtained T_c onset and offset under different pressure is presented in Figure 2(b). Superconductivity is observed in FeSe below found to increase by application of pressure in contrast to conventional superconductors and reaches to 20.2 K under 1.98 GPa, accompanying broadening of transition. However the offset transition temperature which is 6.2 K in ambient pressure is almost doubled under hydrostatic pressure of 1.98 GPa. Due to broadening of transition upon application of higher pressures the onset and offset temperature deviates. This rather broad transition width indicates the inhomogeneous nature of the sample. The change of T_c under pressure may be related to softness of Fe₂Se₂ planes as shown in Figure 1(b), which have significant influence on electronic properties. It

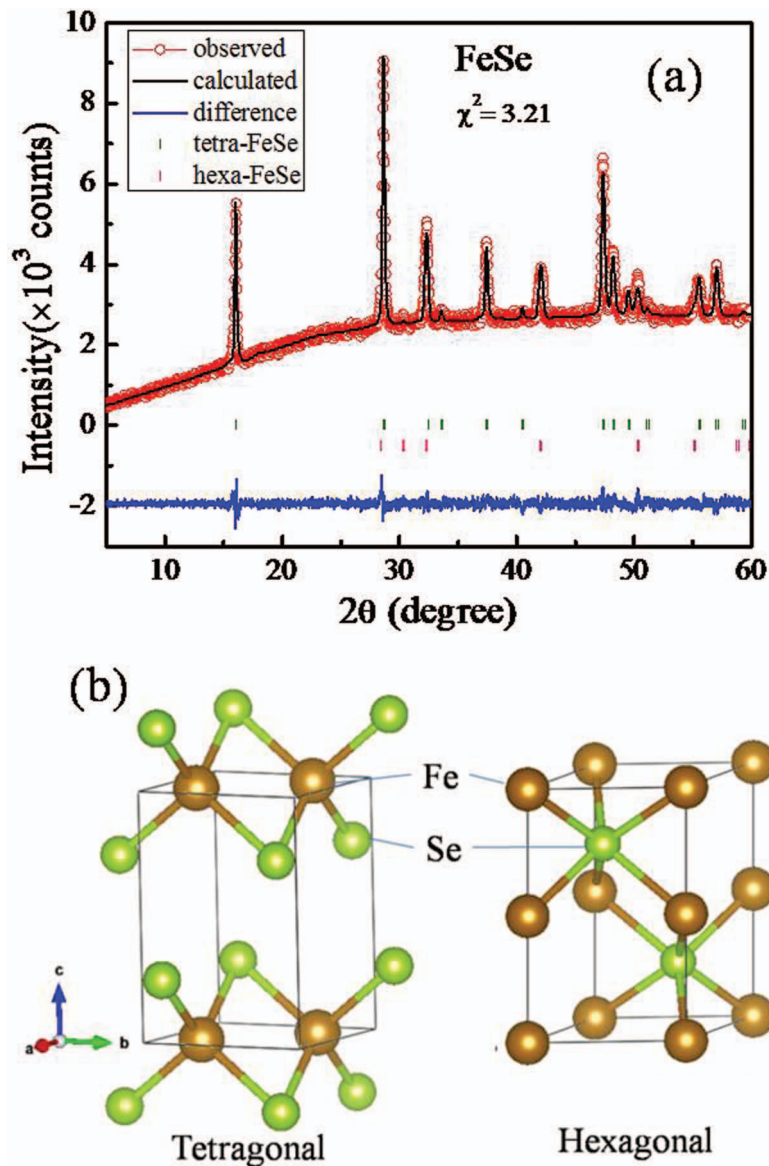


FIG. 1. (a) Rietveld refinement for powder x-ray diffraction pattern of FeSe in both tetragonal P4/nmm (green bars) and hexagonal P63/mmc (red bars) phases. (b) Crystal structures for both tetragonal and hexagonal phases.

was noticed based on Mossbauer study by Medvedev *et al.* that local environment around Fe is not much affected by pressure.¹⁸

To further study the electronic transport of FeSe under hydrostatic pressures of 0 GPa and 1.98 GPa, resistivity measurements as a function of temperature were carried out under different applied magnetic fields parallel to the current flow direction. Under ambient pressure (0 GPa) magnetotransport measurements is displayed in Figure 3(a) where resistive transition shifts towards lower temperature upon increasing magnetic field with broadening of transition width. Under 1.98 GPa hydrostatic pressure the resistive transition as displayed in Figure 3(b) does not show an increase in transition width with applied magnetic field in contrast to 0 GPa case. This suggests a clear enhancement of flux pinning in FeSe sample under applied pressure.

Shown in Figure 4 is the temperature dependence of fields $H_{c2}^{90\%}$ and $H_{c2}^{10\%}$ evaluated at 90% and 10% of normal state resistivity, ρ_N . The temperature dependence of upper critical fields are very close to linear allowing us to use the Werthamer Helfand and Hohenberg (WHH) model for the

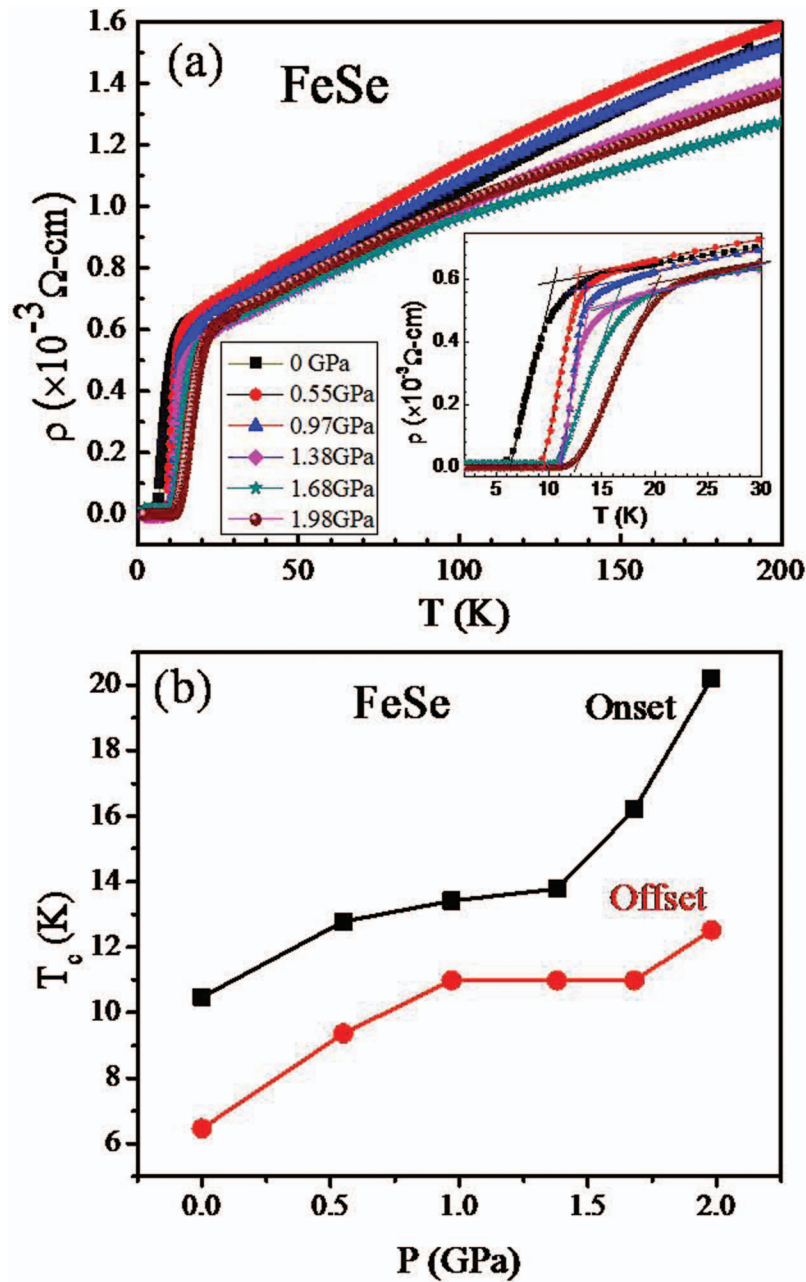


FIG. 2. (a) Resistivity of FeSe under different hydrostatic pressures as function of temperature with a clear increase in superconducting transition is shown for clarity as inset. (b) Variation of onset and offset transition temperature with applied hydrostatic pressure.

determination of uppercritical field $H_{c2}(0)$ by $H_{c2}(0) = -0.69 T_c \left(\frac{dH_{c2}}{dT} \right)_{T=T_c}$ relation within weak coupling BCS superconductors.²⁰ The rate of decrease of upper critical field, dH_{c2}/dT , using the 90% ρ_N and 10% ρ_N criteria are -3.69 Tesla/K and -3.00 Tesla/K which decrease to -3.40 Tesla/K and -2.91 Tesla/K respectively under hydrostatic pressure of 1.98 GPa. We estimated the upper critical field using WHH formula which surprisingly increases from 26.7 Tesla to 47.5 Tesla under 1.98 GPa pressure. The upper critical field value under ambient pressure is close to earlier reported values.^{16–19} In a very recent work by Miyoshi *et al.*, on single crystal of FeSe specimen, the increase in T_c upto 34 K under hydrostatic pressure of 7 GPa is observed, however the does

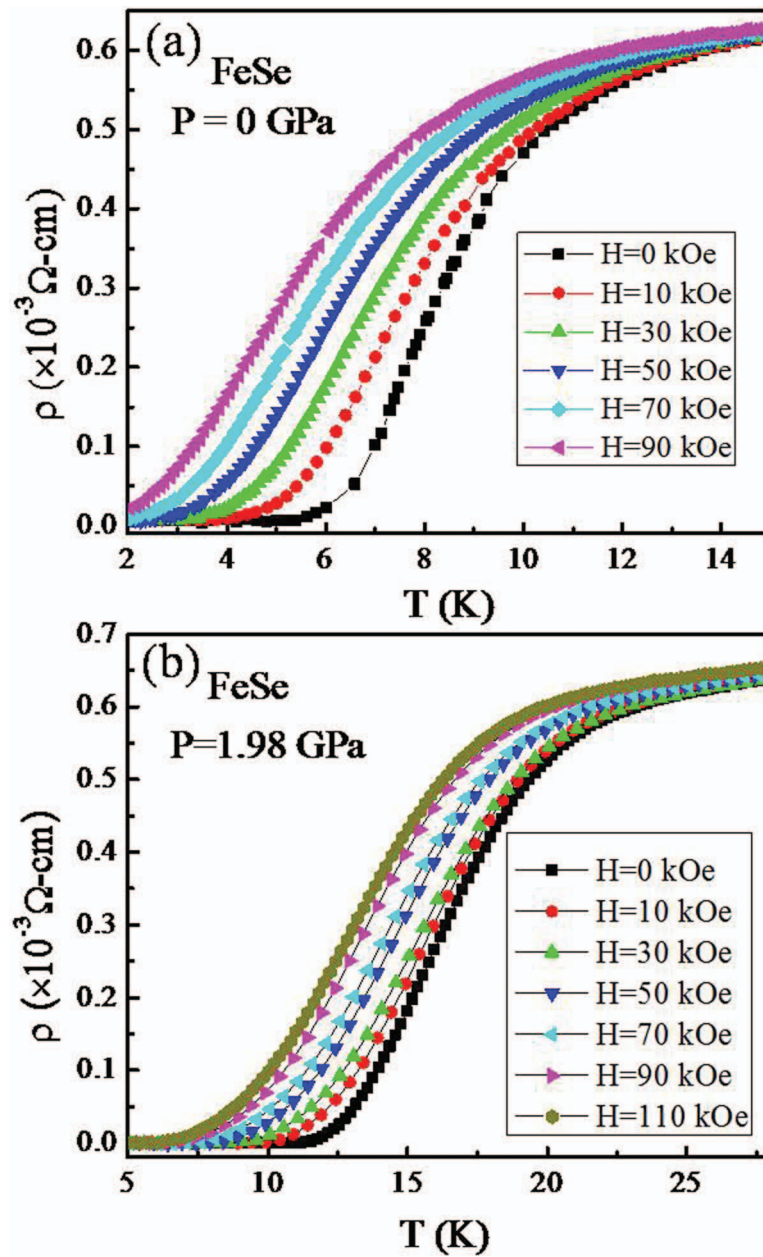


FIG. 3. Resistivity (ρ) versus temperature (T) at different magnetic fields under (a) atmospheric pressure, $P = 0.0$ GPa and (b) $P = 1.98$ GPa.

not comment on upper critical field and pinning.²¹ The superconducting coherence length $\xi(0)$ at absolute zero can be estimated from $H_{c2}(0) = \frac{\Phi_0}{2\pi\xi^2(0)}$ based on Ginzburg-Landau theory where flux quantization $\Phi_0 = 2.07 \times 10^{-15}$ Web. The estimated coherence length decreases under 1.98 GPa pressure from 3.51 nm to 2.63 nm indicating enhanced robustness. The decreased coherence length and enhanced pinning by hydrostatic pressure in conjugation with increased T_c may lead a way to design superconductors for potential applications.

In summary we have studied superconducting properties of synthesized FeSe via the magnetotransport measurements under different hydrostatic pressures. XRD analysis indicates that the superconducting tetragonal FeSe have Fe-vacancy with a minor hexagonal phase. Under hydrostatic pressure of just 1.98 GPa, the superconducting onset temperature increase from 10.4 K to 20.2 K

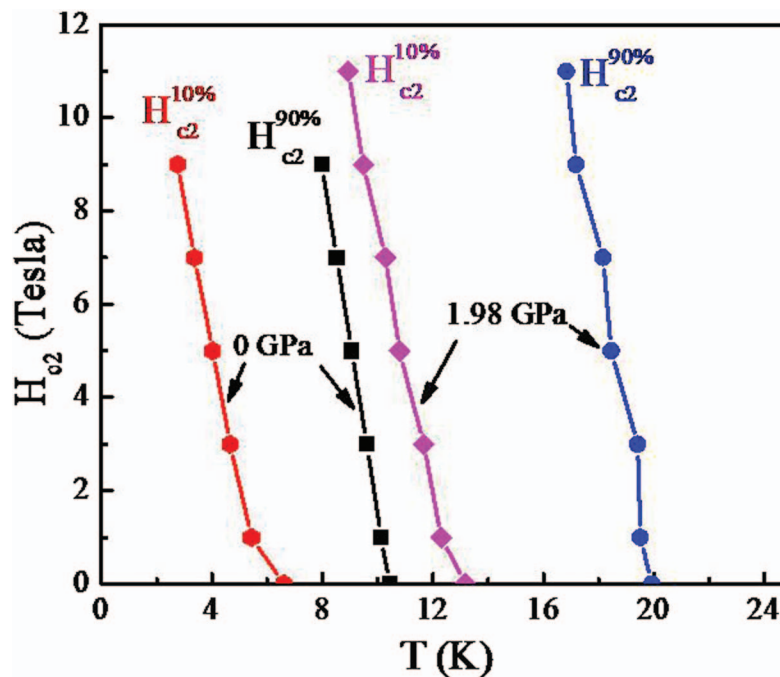


FIG. 4. Variation of upper critical field, $H_{c2}(T)$, with temperature as evaluated from 90% and 10% of normal state resistivity T_c (50% ρ_n) under atmospheric and 1.98 GPa hydrostatic pressure. Solid lines are the guide to eyes.

with broadening in transition. From the magnetotransport measurements the estimated upper critical field is 26.7 Tesla which increased to 47.5 Tesla under increased pressure which in turn results in decreased coherence length in conjugation to enhanced flux pinning.

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